

Determination of the 1-D distribution of electrical conductivity in Earth's mantle from *Swarm* satellite data

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We present an inversion scheme to recover the (1-D) depth profile of mantle conductivity from satellite magnetic data, which takes into account 3-D effects arising from the distribution of oceans and continents. The scheme is based on an iterative inversion of C -responses, which are estimated from time series of the dominating external (inducing) and internal (induced) spherical harmonic coefficients of the magnetic potential due to a magnetospheric source. These time series will be available as a *Swarm* Level-2 data product. We verify our approach by using synthetic, but realistic time series obtained by simulating induction due to a realistic magnetospheric source in a 3-D “target” conductivity model of the Earth. This model contains not only a laterally heterogeneous layer representing oceans and continents, but also 3-D inhomogeneities in the mantle. The inversion for mantle conductivity is initiated with a uniform conductivity model. Convergence is reached within a few iterations. The recovered model agrees well with the laterally averaged target model, although the latter comprises large jumps in conductivity. Our 1-D inversion scheme is therefore ready to process *Swarm* data.

Key words: Electromagnetic induction, 1-D inversion, mantle conductivity, C -responses.

1. Introduction

Temporal variations of Earth's magnetic field have long been used to infer global 1-D conductivity profiles of Earth's mantle, mostly from continental geomagnetic observatories (e.g. Schmucker, 1985; Schultz and Larsen, 1987; Olsen, 1998). In contrast to data from observatories, which are sparse and irregularly distributed (with only few in oceanic regions), data of uniform quality with a good spatial coverage can be obtained from low-Earth-orbit (LEO) platforms. This enables the derivation of a globally-averaged conductivity profile that is not biased towards continental regions.

However, analysis of satellite data is more challenging compared to analysis of observatory data for two reasons: First, LEO satellites move typically with a speed of 7–8 km/s and thus measure a mixture of temporal and spatial changes of the magnetic field. Second, satellites pass over both continents and oceans, and therefore the magnetic satellite data are affected by induction in the oceans (cf. Tarits and Grammatica, 2000; Everett *et al.*, 2003; Kuvshinov and Olsen, 2005) in a complicated way. In spite of these difficulties, a number of attempts has been made to probe mantle conductivity from space (e.g. Didwall, 1984; Olsen, 1999; Olsen *et al.*, 2002; Constable and Constable, 2004; Velínský *et al.*, 2006). Most recently, Kuvshinov and Olsen (2006) derived a global model of mantle conductivity from five years of CHAMP, Ørsted and SAC-C magnetic data and demonstrated the necessity of taking into ac-

count the contributions of a laterally heterogeneous surface shell representing the distribution of oceans and continents. The authors accounted for the ocean effect by correcting the magnetic field at orbit altitudes.

As in the paper by Kuvshinov and Olsen (2006), we present in this paper a methodology to determine the 1-D conductivity profile of Earth's mantle by inverting global C -responses. These are estimated from time series of the dominating external (inducing) and internal (induced) spherical harmonic expansion (SHE) coefficients of the magnetic potential that describes the signals of magnetospheric origin. Both time series will be available as *Swarm* Level-2 data product MMA_SHA_2_, determined by the Comprehensive Inversion chain (CI, Sabaka *et al.*, 2013) of the *Swarm* Level-2 Processing Facility SCARF (Olsen *et al.*, 2013). The CI aims to separate magnetic contributions from various sources (originating in the core, lithosphere, ionosphere and magnetosphere) in the form of corresponding SHE coefficients. As Kuvshinov and Olsen (2006), we also correct for the ocean effect, but our correction scheme differs from that presented by the authors in two aspects. First, our scheme is iterative, i.e. a correction for the ocean effect is applied multiple times. Second, we do not make the detour of predicting the magnetic field at orbit altitudes, but directly correct the estimated C -responses.

In Section 2 of this paper, we outline the inversion algorithm and describe how we account for the ocean effect. Section 3 presents results of a model study. We summarize our work in Section 4.

2. Inversion Algorithm

In this section, we outline the succession of processing steps that forms the inversion scheme. A summary is pre-

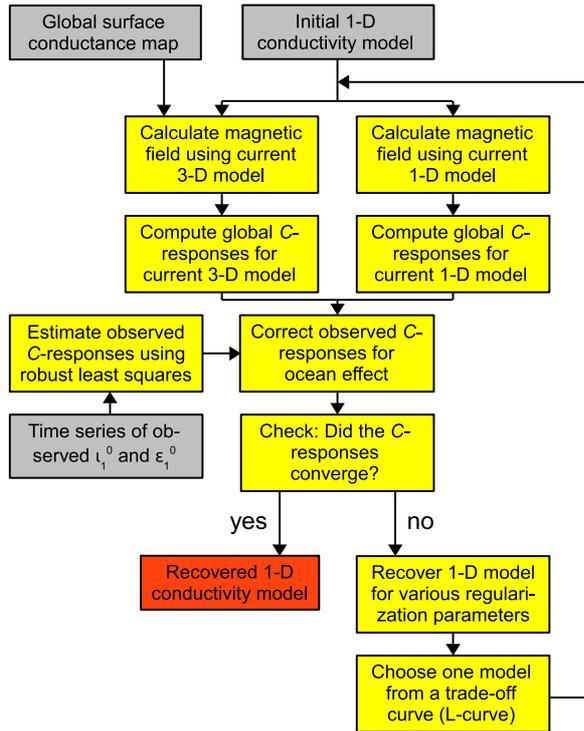


Fig. 1. Scheme of the processing steps (yellow boxes) that constitute the iterative 1-D inversion. Inputs are marked by grey boxes, outputs by orange boxes. Note that by the term “3-D model”, we denote the 1-D model plus a laterally heterogeneous surface shell.

sented in Fig. 1. Note that our inversion scheme is iterative. The iterative structure is due to the applied correction for the ocean effect (arising from a laterally heterogeneous surface shell). Without accounting for this effect, only the estimate of global C -responses from the observed data (Subsection 2.1) and the subsequent recovery of a conductivity model (Subsection 2.3) would be necessary. However, the C -responses estimated in such a way would be biased at short periods, thus also biasing the recovered conductivity model (mainly at shallow depths, cf. Kuvshinov and Olsen, 2006).

2.1 Estimation of global C -responses

In the source-free region above the conducting Earth, the magnetic field due to a magnetospheric source can be represented as gradient of a scalar potential, $\mathbf{B} = -\nabla V$, which is obtained by solving Laplace’s equation. The potential V can be expanded into external and internal spherical harmonic sources. We assume a ring current geometry for the external part of V , described by a single spherical harmonic coefficient, $\varepsilon_1^0(\omega)$. In a 1-D Earth, a source with this geometry only induces one internal coefficient, $\iota_1^0(\omega)$. Thus, V can (in frequency domain) be written as

$$V(r, \vartheta, \omega) = a \left\{ \varepsilon_1^0(\omega) \left(\frac{r}{a}\right) + \iota_1^0(\omega) \left(\frac{a}{r}\right)^2 \right\} \cos \vartheta, \quad (1)$$

where ϑ denotes colatitude, r denotes distance from Earth’s centre, a is Earth’s mean radius and ω is angular frequency. The 1-D electromagnetic transfer function of de-

gree 1, $Q_1(\omega)$, is defined by the relation

$$\iota_1^0(\omega) = Q_1(\omega) \varepsilon_1^0(\omega). \quad (2)$$

Time series of ε_1^0 and ι_1^0 , provided by the CI, are the inputs for our inversion. We use the section averaging approach (Olsen, 1998) and a robust statistical procedure involving iteratively re-weighted least squares (Aster *et al.*, 2005) to estimate the transfer function $Q_1^{\text{obs}}(\omega)$ (and the corresponding uncertainties $\delta Q_1^{\text{obs}}(\omega)$) from ε_1^0 and ι_1^0 at a set of (logarithmically spaced) frequencies ω . The 1-D Q -response is transformed to the global C -response by means of

$$C^{\text{obs}}(\omega) = \frac{a}{2} \frac{1 - 2Q_1^{\text{obs}}(\omega)}{1 + Q_1^{\text{obs}}(\omega)}, \quad (3)$$

with corresponding errors

$$\delta C^{\text{obs}}(\omega) = \frac{3a}{2} \frac{1}{|1 + Q_1^{\text{obs}}(\omega)|^2} \delta Q_1^{\text{obs}}(\omega), \quad (4)$$

which follows from the error propagation law.

2.2 Correction of estimated C -responses

Within each iteration of the inversion scheme, we simulate induction in a 1-D model and a 3-D model, yielding the synthetic global C -responses $C^{1-D}(\omega)$ and $C^{3-D}(\omega)$, respectively (cf. Fig. 1). The 1-D model consists of the 1-D conductivity structure recovered in the previous iteration of the inversion. With the term “3-D model”, we denote this 1-D model plus a laterally heterogeneous surface shell. The calculation of synthetic C -responses for a X -D model, where X -D refers to either 1-D or 3-D, involves three steps:

- 1) Calculation of $\mathbf{B}^{\text{X-D,unit}}(r = a, \vartheta, \varphi, \omega)$, i.e. the magnetic field at Earth’s surface due to a unit amplitude magnetospheric ring current source at a set of frequencies ω . Note that the calculated magnetic field only varies in longitude φ if the model contains 3-D heterogeneities.
- 2) Recovery of the transfer function $Q_1^{\text{X-D}}(\omega)$ by spherical harmonic analysis of $B_r^{\text{X-D,unit}}(r = a, \vartheta, \varphi, \omega)$ using the formula

$$Q_1^{\text{X-D}}(\omega) = \frac{3}{8\pi} \int_S (B_r^{\text{X-D,unit}}(r = a, \vartheta, \varphi, \omega) + \cos \vartheta) \cos \vartheta ds, \quad (5)$$

where $ds = \sin \vartheta d\vartheta d\varphi$. Equation (5) follows from $\mathbf{B} = -\nabla V$ with V given by Eq. (1), the definition of the Q -response (2) and $\varepsilon_1^0(\omega) = 1$.

- 3) Transformation of $Q_1^{\text{X-D}}(\omega)$ to $C^{\text{X-D}}(\omega)$ using Eq. (3).

To compute the magnetic field in a 3-D conductivity model (item 1 of the above list), we use a 3-D contracting integral equation solver, extensively described in Kuvshinov and Semenov (2012).

We correct the global C -responses estimated from observed data, $C^{\text{obs}}(\omega)$ (cf. Subsection 2.1), for the ocean effect with the formula

$$C^{\text{corr}}(\omega) = C^{\text{obs}}(\omega) + C^{1-D}(\omega) - C^{3-D}(\omega). \quad (6)$$

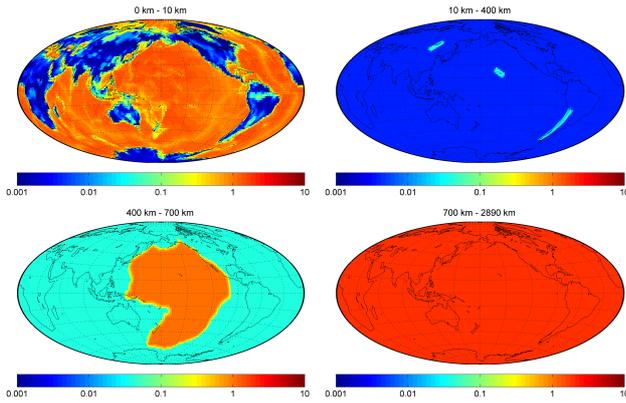


Fig. 2. Target conductivity model used in our model study, units are in S/m. Note that the conductivity of the top layer has been obtained by scaling the surface conductance map to a thickness of 10 km.

This correction diminishes 3-D effects (arising from the laterally heterogeneous surface shell) in the data. If our 3-D conductivity model coincides with the true conductivity structure of the Earth, $C^{3-D}(\omega)$ and $C^{\text{obs}}(\omega)$ cancel out except for measurement errors, and $C^{\text{corr}}(\omega)$ is then simply given by $C^{1-D}(\omega)$. This logic is also applied to decide when to stop iterating. Convergence of the iterative scheme is reached if the weighted RMS of $C^{\text{obs}}(\omega)$ and $C^{3-D}(\omega)$ falls below a prescribed threshold ϵ , i.e. if

$$\text{RMS} = \sqrt{\frac{1}{N_\omega} \sum_{j=1}^{N_\omega} \frac{|C^{\text{obs}}(\omega_j) - C^{3-D}(\omega_j)|^2}{\delta C^{\text{obs}}(\omega_j)^2}} \leq \epsilon, \quad (7)$$

where N_ω is the number of periods at which responses were estimated. If, on the other hand, the RMS is larger than ϵ , the corrected C -responses are inverted for a new 1-D conductivity model, and a new iteration is initiated (cf. Fig. 1).

2.3 Derivation of the 1-D conductivity model

We derive the 1-D conductivity model from the corrected C -responses $C^{\text{corr}}(\omega)$ by using the quasi-Newton algorithm of Byrd *et al.* (1995). The inversion is stabilized by minimizing the first derivative of $\log(\text{conductivity})$ with respect to $\log(\text{depth})$. Inversion is performed for several regularization parameters (i.e. several degrees of smoothing). The solution is picked from a trade-off curve (L-curve), which relates data misfit and model complexity (Hansen, 1992).

3. Model Study

In order to test the performance of our inversion scheme, we generate synthetic data (i.e. time series of induced coefficients) in a test 3-D conductivity model, hereinafter referred to as “target model”, and afterwards recover the 1-D conductivity structure of the target model (i.e. its laterally averaged conductivity) from the data. We introduce the target conductivity model (Subsection 3.1), describe how we generate the test data (Subsection 3.2) and finally present the results of our model study (Subsection 3.3).

3.1 Target conductivity model

Figure 2 shows the target conductivity model. It consists of a thin surface shell of laterally varying conductance and a layered model, which contains different con-

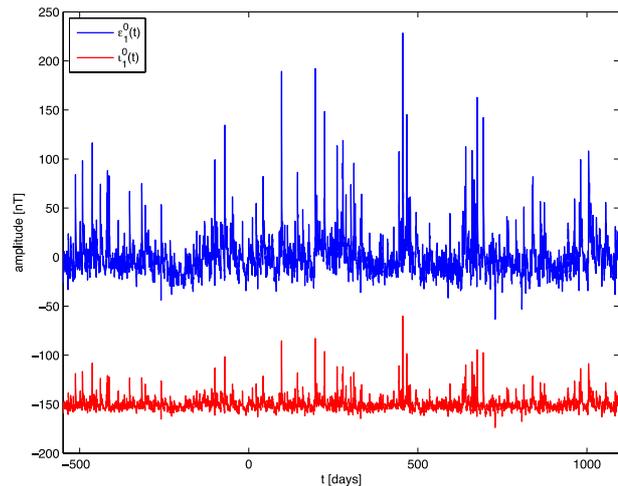


Fig. 3. Input data of the model study—time series of ε_1^0 and t_1^0 (in nT). The time (in days) is relative to January 1, 2000. Note that the time series of t_1^0 has been shifted by 150 nT for clarity.

ductivity anomalies, underneath. The surface shell is scaled to a thickness of 10 km. The conductivity anomalies in deeper regions are introduced in order to simulate that the true Earth most probably also contains 3-D mantle heterogeneities. Also note that the same target model has been used in the associated 3-D mantle conductivity inversion studies (cf. companion papers by Püthe and Kuvshinov, 2013; Velínský, 2013).

3.2 Generation of the test data set

Hourly mean time series of external SHE coefficients in a geomagnetic dipole coordinate system (up to degree 3 and order 1) have been derived by analysis of 4.5 years of observatory data (July 1998–December 2002), details of the derivation are given in Olsen *et al.* (2005). Time series of internal SHE coefficients are synthesized by simulating induction in the target conductivity model (Subsection 3.1), cf. companion paper by Püthe and Kuvshinov (2013).

We use the time series of external and internal coefficients to predict the magnetospheric field at orbit altitudes (with a sampling frequency of 1 Hz). Adding the contributions due to different sources (core, lithosphere and ionosphere) yields the magnetic field at orbit altitudes, which is then analyzed by the CI. The external and internal SHE coefficients of the magnetic potential due to magnetospheric sources recovered by the CI constitute a realistic test data set for our inversion scheme. A more detailed description of the generation of the test data set is provided in Olsen *et al.* (2013).

3.3 Inversion results

Recovered time series of ε_1^0 and t_1^0 are provided by the CI with a sampling rate of 1.5 hours (Level-2 data product MMA_SHA_2_, cf. Sabaka *et al.*, 2013). The time series are depicted in Fig. 3. With these data, we estimate C -responses at 23 logarithmically spaced periods between 14 hours and 83 days. We invert the C -responses to recover the 1-D mantle conductivity at depths between 10 km and the core-mantle boundary at 2890 km. The inversion domain is stratified into in total 44 layers with thicknesses of 50 km

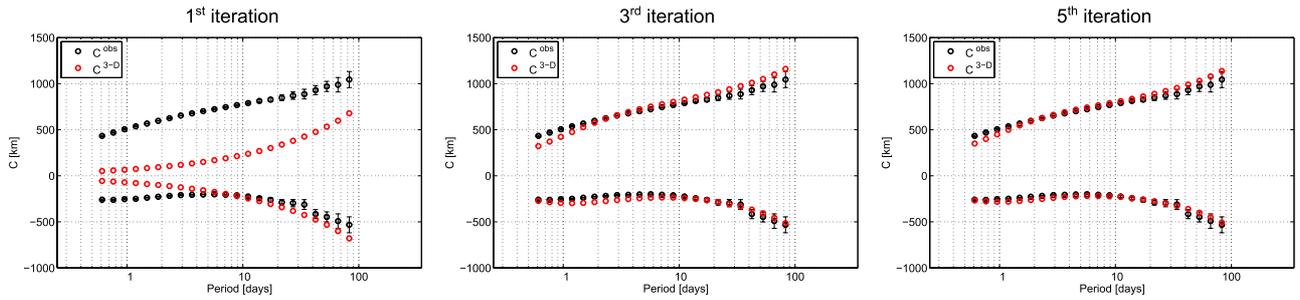


Fig. 4. Demonstration of the convergence of observed global C -responses $C^{\text{obs}}(\omega)$ and C -responses calculated in the current 3-D model, $C^{3\text{-D}}(\omega)$, within five iterations of the inversion scheme. Note that the series with positive values correspond to $\text{Re } C$, those with negative values to $\text{Im } C$.

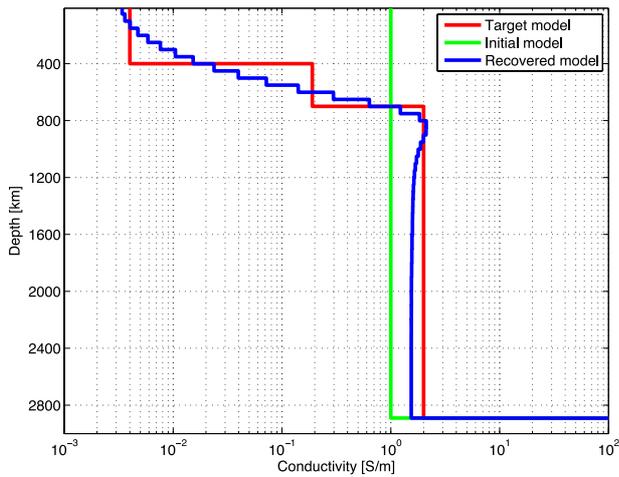


Fig. 5. Recovered conductivity model (blue) in comparison to the uniform initial model (green) and the laterally averaged target conductivity model (red; compare with Fig. 2). Note that the conductivity of the core is fixed to 10^5 S/m.

(at depths below 1500 km) and 100 km (at depths greater than 1500 km), respectively. As initial (starting) model, a uniform mantle with conductivity of 1 S/m is prescribed.

For a chosen threshold value $\epsilon = 5$ (cf. Eq. (7)), the inversion scheme converged after 5 iterations. Figure 4 shows the convergence of the C -responses. The recovered conductivity model in comparison to the (laterally averaged) target model is shown in Fig. 5. Lateral averaged conductivity here denotes the arithmetic mean of the conductivity of all cells in the respective layer. The results indicate that the inversion scheme is able to accurately recover mantle conductivity at all depths. Although the conductivity of the initial model is very different from the target conductivity structure, the final model agrees well with the target model. Due to the applied smoothing, the recovered model does not comprise the large jumps in conductivity that are apparent in the target model at depths of 400 km and 700 km. Such large jumps in conductivity are, however, not likely for true Earth.

4. Conclusions

We have presented a scheme to invert satellite magnetic data for a global depth profile of mantle conductivity. The scheme is based on the inversion of C -responses and comprises a correction for the ocean effect. In spite of the it-

erative architecture and a number of processing steps, final results can be obtained within a few hours. The repeated calculation of the magnetic field in a 3-D model is the most expensive step in terms of computational cost.

The algorithm has been tested by simulating induction due to a realistic magnetospheric source in a realistic 3-D (target) conductivity model and recovering the 1-D conductivity structure of this model (i.e. its laterally averaged conductivity) from the synthetic data. In spite of large conductivity jumps and a number of 3-D heterogeneities in the target model, an excellent recovery of the 1-D conductivity of Earth's mantle has been achieved from crustal depths to the core-mantle boundary. The algorithm has thus proved to be workable and ready to digest *Swarm* data. Moreover, the inversion results provide a useful initial guess for 3-D mantle conductivity studies with *Swarm* data (cf. companion papers by Püthe and Kuvshinov, 2013; Velínský, 2013).

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